

EP 403  
Final Exam 1991

1. (15 marks)

a. (5 marks)

Boron is a good absorber,  $\therefore$  safe to assume negligible scattering (no buildup).

$$\therefore \frac{I(x)}{I(0)} = 0.001 = e^{-\Sigma_a x}$$

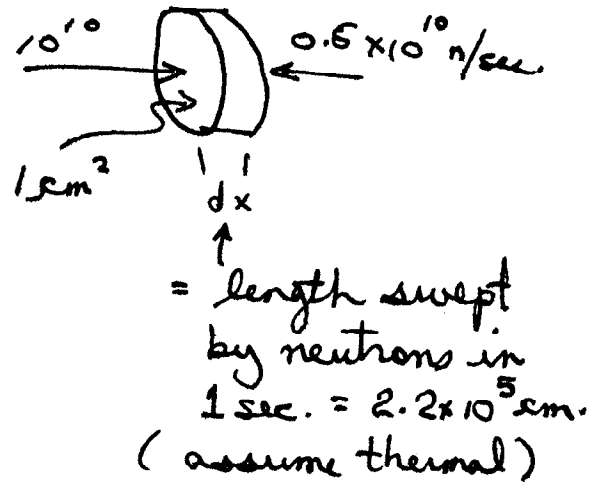
$$\therefore x = -\frac{\ln(0.001)}{\Sigma_a} = \frac{+6.91}{103}$$

$\Sigma_a \uparrow 103 \text{ cm}^{-1}$

$$= \underline{\underline{0.0671 \text{ cm.}}}$$

b. (10 marks)

$$\begin{aligned} \phi &= \frac{\int_V n v dV}{\int_V dV} \\ &= \frac{1.5 \times 10^{10} \times 2.2 \times 10^5}{2.2 \times 10^5} \\ &= 1.5 \times 10^{10} \text{ neutrons} \\ &\quad \text{cm}^2\text{-sec.} \end{aligned}$$



$$\begin{aligned} J &= \int_S n \underline{v} \cdot d\underline{s} = (1.0 \times 10^{10} - 0.5 \times 10^{10}) \cdot 1 \text{ cm}^2 \\ &= 0.5 \times 10^{10} \hat{x} \text{ neutrons/cm}^2\text{-sec.} \\ &\quad \text{(positive x direction)} \end{aligned}$$

2.

(10 marks)  
total

$$\frac{dN}{dt} = -\lambda N - cN + R$$
$$= -(\lambda+c)N + R$$

(5 marks for  
correct rate  
eqn.)

Let  $x(t) = N(t) - R$  ~~eqn.~~  $\cdot a$

$\uparrow$  a constant  
to be determined

$$\therefore \frac{dx}{dt} = -(\lambda+c)[x + aR] + R$$

Let  $a = \frac{1}{\lambda+c}$  to give

$$\frac{dx}{dt} = -(\lambda+c)x$$

$$\Rightarrow x = x_0 e^{-(\lambda+c)t}$$

(3 marks for  
a correct  
sol'n procedure)

where  $x_0 = N(0) - \frac{R}{\lambda+c}$

$$\therefore N(t) = \left[ N(0) - \frac{R}{\lambda+c} \right] e^{-(\lambda+c)t} + \frac{R}{\lambda+c}$$

$$\therefore N(t) = N(0) e^{-(\lambda+c)t} + \frac{R}{\lambda+c} (1 - e^{-(\lambda+c)t})$$

$\uparrow$   
decay + removal

$\uparrow$   
production

(2 marks for  
correct  
solution)

3. (10 marks)

a.  $\phi = \frac{SL}{2D} e^{-x/L}$ ,  $x > 0$

Absorption rate at  $x = \Sigma_a \phi(x)$

Total absorption rate for  $x > 0$ :

$$= \int_0^{\infty} \Sigma_a \phi(x) dx \quad \#/\text{cm} \cdot \frac{1}{\text{cm}^2\text{-sec.}}$$

$$= \frac{\Sigma_a SL}{2D} \int_0^{\infty} e^{-x/L} dx$$

$$= \frac{\Sigma_a SL}{2D} (-L) e^{-x/L} \Big|_0^{\infty}$$

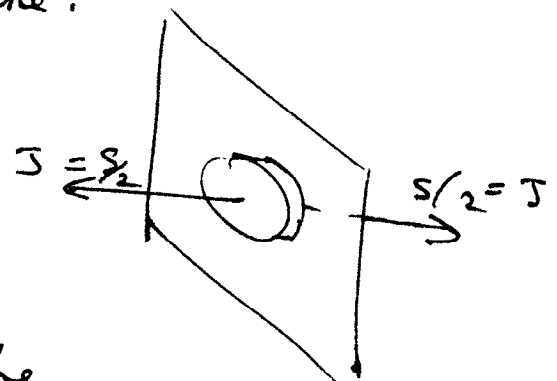
$$= \frac{S}{2} \quad (7 \text{ marks})$$

(4 marks for properly posing integral.)

3 for proper solution.)

b. Current at source plane:

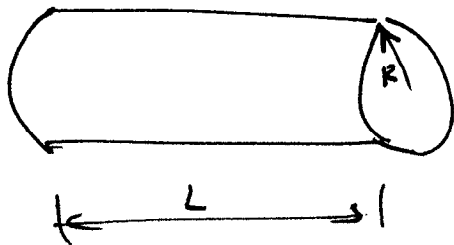
$$= \lim_{x \rightarrow \infty} J(x) = \frac{S}{2}$$



$\therefore$  As you might expect, the current flowing to the right equals exactly the absorption, i.e. sinks = sources in steady state. (3 marks)

4. (10 marks)

$$\nabla^2 \phi + B_g^2 \phi = 0 \Rightarrow B_g^2 = \left(\frac{\pi}{L}\right)^2 + \left(\frac{v_0}{R}\right)^2$$



for criticality  
Volume =  $\pi R^2 L$

For volume to be minimum:  $\frac{\partial \text{Volume}}{\partial R} = \frac{\partial \text{Volume}}{\partial L} = 0$

Since  $L$  &  $R$  are constrained via  $B_g$ :

$$\frac{v_0}{R} = \sqrt{B_g^2 - \left(\frac{\pi}{L}\right)^2} \Rightarrow R = \frac{v_0}{\sqrt{B_g^2 - \left(\frac{\pi}{L}\right)^2}}$$

$$\therefore \text{Volume} = \pi R^2 L = \frac{\pi v_0^2 L}{\left[B_g^2 - \left(\frac{\pi}{L}\right)^2\right]}$$

$$\therefore \frac{\partial \text{Volume}}{\partial L} = \frac{\pi v_0^2}{\left[B_g^2 - \left(\frac{\pi}{L}\right)^2\right]} \Rightarrow \frac{\pi v_0^2 L \left(+ 2\pi^2/L^3\right)}{\left[B_g^2 - \left(\frac{\pi}{L}\right)^2\right]^2} = 0$$

$$\therefore 1 = \frac{2\pi^2}{L^2 \left[B_g^2 - \left(\frac{\pi}{L}\right)^2\right]} = 0$$

$$\therefore \left(\frac{v_0}{R}\right)^2 = 2\left(\frac{\pi^2}{L^2}\right) \Rightarrow \text{Radial buckling} = 2 \times \text{axial buckling}$$

$B_g^2$  radial

$B_g^2$  axial

(6 marks for formulation of problem, 4 for solution)

5. (15 marks)

a. (4 marks)

$$-D_1^c \nabla^2 \phi_1^c + \sum_{R_1}^c \phi_1^c = \frac{1}{K} \left[ \nu_1 \sum_{F_1}^c \phi_1^c + \nu_2 \sum_{F_2}^c \phi_2^c \right]$$

$$-D_2^c \nabla^2 \phi_2^c + \sum_a^c \phi_2^c = \sum_{S_{12}}^c \phi_1^c$$

$$-D_1^R \nabla^2 \phi_1^R + \sum_{R_1}^R \phi_1^R = 0$$

$$-D_2^R \nabla^2 \phi_2^R + \sum_a^R \phi_2^R = \sum_{S_{12}}^R \phi_1^R$$

b. (4 marks)

$$\phi_i^R \left( \frac{b}{2} \right) = 0 \quad i = 1, 2$$

$$\phi_i^c \left( \frac{a}{2} \right) = \phi_i^R \left( \frac{a}{2} \right)$$

$$\mathcal{J}_i^c \left( \frac{a}{2} \right) = \mathcal{J}_i^R \left( \frac{a}{2} \right)$$

$$\left( \text{ie } D_i^c \nabla \phi_i^c \left( \frac{a}{2} \right) = -D_i^R \nabla \phi_i^R \left( \frac{a}{2} \right) \right)$$

$$\text{Symmetry or } \nabla \phi_i^c(0) = 0$$

c. (4 marks)

Try solutions of form:  $\phi_i^* = A_i \cos \mu_i r + C_i \cosh \lambda_i r$  (or whatever, Bessel?)  
✓ subst. into eqn's in (a). This yields 4 eqn's that should allow the calc. of the buckling

coefficients,  $\mu_i$  &  $\lambda_i$ .

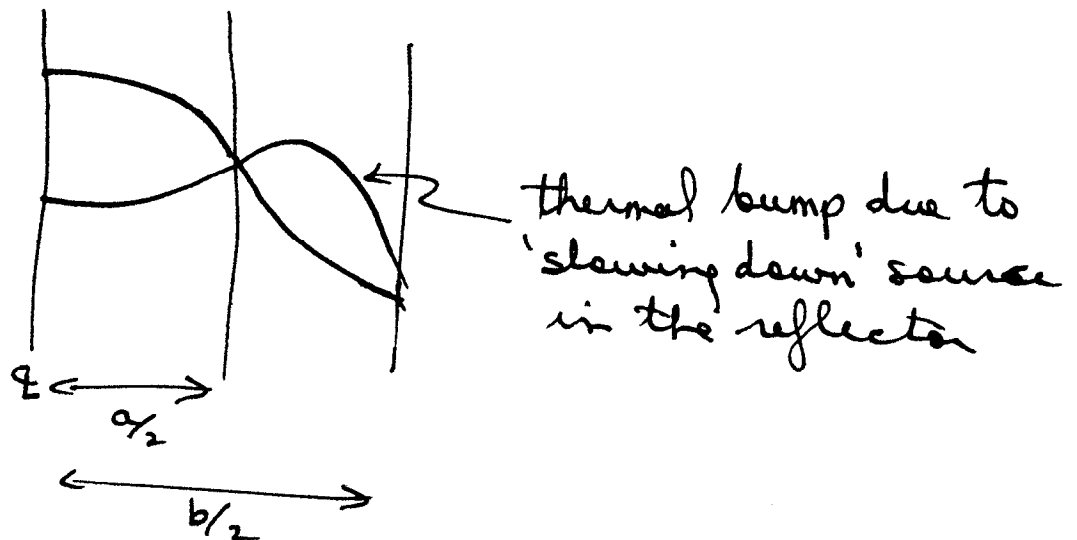
Next apply the B.C.'s. The condition

$\phi_i^R(a/2 + b) = 0$  should allow you to eliminate one unknown  $A_i^R$  or  $C_i^R$  for  $i=1,2$ .

Symmetry does the same for the core flux.

This leaves 4 unknowns & 4 interface B.C. equations that can be solved (theoretically) for the 4 unknown. These 4 eqn's form the criticality condition.

d) (3 marks)



6. (15 marks)

a. Tightly Coupled: scattering down to nearest neighbour only (5 marks)

$\Rightarrow$   $\left( \parallel \right)$

No upscatter: can't scatter up

$\Rightarrow$   $\triangle$

$$\sum_{s g' g} = 0 \text{ for } g' > g$$

$$\therefore \sum \sum_{s g g'} \phi_{g'} = \sum_{g'=1}^{g-1} \sum_{s g' g} \phi_g + \sum_{s g g} \phi_g$$

$\Rightarrow$  lump with  $\Sigma_{Rg}$

$$\sum_{g'=1}^G \sum_{s g' g} \phi_{g'} = \sum_{s g-1, g} \phi_g + \sum_{s g g} \phi_g$$

b. As per text. (5 marks)

$$\text{eg } \Sigma_{ag} = \frac{\int_{E_g}^{E_{g-1}} \Sigma_a(E) \phi(E) dE}{\int_{E_g}^{E_{g-1}} \phi dE}$$

$$\text{Thus } \Sigma_a = \sum_{g=1}^G \Sigma_{ag} \phi_g / \sum_{g=1}^G \phi_g = \phi$$

yields:

$$\frac{\partial \phi}{\partial t} = \nabla \cdot D \nabla \phi - \Sigma_a \phi - \Sigma_s \phi_g + \Sigma_s \phi_g$$
$$+ \nu \Sigma_f \phi + \sum_{i=1}^N \lambda_i C_i + S_{ext}$$

$$\frac{\partial C_i}{\partial t} = -\lambda_i C_i + \beta_i \nu \Sigma_f \phi$$



c. (5 marks)

If  $t_{1/2}$  small, then  $\lambda_i$  large

$\therefore$  if  $C_i$  is significant,  $\frac{\partial C_i}{\partial t}$  is large & -ve.

$\therefore$  precursors decay rapidly away.

$\therefore C_i \approx 0$  &  $\frac{\partial C_i}{\partial t} \approx 0$

$$\therefore 0 = -\lambda_i C_i + \sum_{g=1}^G \beta_{ig} \nu_g \Sigma_{fg} \phi_g$$

$$\therefore \sum_{i=1}^N \lambda_i C_i = \sum_{g=1}^G \nu_g \Sigma_{fg} \phi_g \underbrace{\sum_{i=1}^N \beta_{ig}}_{\equiv \beta_g}$$

$\therefore$  term in flux eqn simplifies to:

$$\chi_g \sum_{g'=1}^G \nu_{g'} (1 - \beta_{g'}) \Sigma_{fg'} \phi_{g'} + \sum_{i=1}^N \lambda_i C_i$$

these 2 terms cancel.

& drop precursor eqn since  $C_i \approx 0$

7. (10 marks)

From the notes we have:

Coolant

$$T_{\text{fluid}}(z) = T_{\text{inlet}} + \frac{q'_0 H}{\pi C_p W} \left[ \sin\left(\frac{\pi z}{H}\right) + 1 \right]$$

$$z \in \left(-\frac{H}{2}, \frac{H}{2}\right)$$

$$q'_0 = q'_0 \cos\left(\frac{\pi z}{H}\right)$$

where  $z \in \left(-\frac{H}{2}, \frac{H}{2}\right)$

Fuel

$$T_{\xi} = T_{\text{fluid}}(z) + \frac{q'_0(z)}{2\pi r_f} \left[ \frac{r_f}{2\bar{k}_f} + \frac{1}{h_G} + \frac{t_c + t_G}{k_c} + \frac{r_f}{h_s(r_f + t_c + t_G)} \right]$$

Since no sheath,  $t_c = t_G = 0$ ,  $h_G = \infty$

$$\therefore T_{\xi}(z) = T_{\text{fluid}}(z) + \frac{q'_0(z)}{2\pi r_f} \left[ \frac{r_f}{2\bar{k}_f} + \frac{1}{h_s} \right]$$

$$= T_{\text{inlet}} + \frac{q'_0 H}{\pi C_p W} \left[ \sin\left(\frac{\pi z}{H}\right) + 1 \right] + \frac{q'_0 \cos\left(\frac{\pi z}{H}\right)}{2\pi r_f} \left[ \frac{r_f}{2\bar{k}_f} + \frac{1}{h_s} \right]$$