

Chapter 7 The States of a System

- Any <sup>practical</sup> system can be in a large number of possible states. The number is large but finite.
- The number increases rapidly with the number of particles in the system and the number of possible states per particle.
- we talked about ways to deal with these states:

Binomial:  $P_N(n) = \frac{N!}{n!(N-n)!} p^n q^{N-n}$   
 (  $\log(m!) \approx m \ln m - m + \frac{1}{2} (2\pi m)$  )

Poisson:  $P_N(n) \approx \frac{(\bar{n})^n}{n!} e^{-\bar{n}}$ ,  $p \ll 1$   
 $\bar{n} = pN$

Gaussian:  $P(n) = \frac{1}{\sqrt{2\pi}\sigma} e^{-\frac{(n-\bar{n})^2}{2\sigma^2}}$

$\sigma^2 = (n-\bar{n})^2 = Npq$ ,  $\bar{n} = pN$   
 $(n-\bar{n}) \ll \sigma^2$

- Quantum mechanics assures us that the number of states is finite (though very large)

## A. Equilibrium

- An isolated system is "in equilibrium" when the probabilities of it being in the various possible states do not vary in time.
- each system has a characteristic relaxation time for each process.
  - heat transfer
  - particle redistribution due to mechanical motion
  - diffusion
  - ...
- Thermodynamics is really Thermo "statics" or thermo "quasi-statics".  
 Changes on a system level are considered to be slow w.r.t. relaxation times
  - i.e. quasi-static.

## B. The Fundamental Postulate

An isolated system in equilibrium is equally likely to be in any of its accessible states.

ie  $P = \frac{1}{\Omega_0}$

↖ # of possible states of system 0

If  $\Omega_i \equiv$  number of states in subset  $i$

then  $P_i =$  prob. of system being in subset  $i$  state

$$= \frac{\Omega_i}{\Omega_0}$$

ie life is a game (lottery)

Example:

3 coins flipped (Heads or Tails)

8 Possible outcomes:

- |   |       |     |  |
|---|-------|-----|--|
| 1 | T T T |     |  |
| 2 | T T H |     |  |
| 3 | T H T |     |  |
| 4 | T H H |     |  |
| 5 | H T T | * X |  |
| 6 | H T H | * ✓ |  |
| 7 | H H T | * ✓ |  |
| 8 | H H H | * X |  |

What is prob. of getting 2 heads out of 3 given the first flip is a head? (Call this outcome 1)?

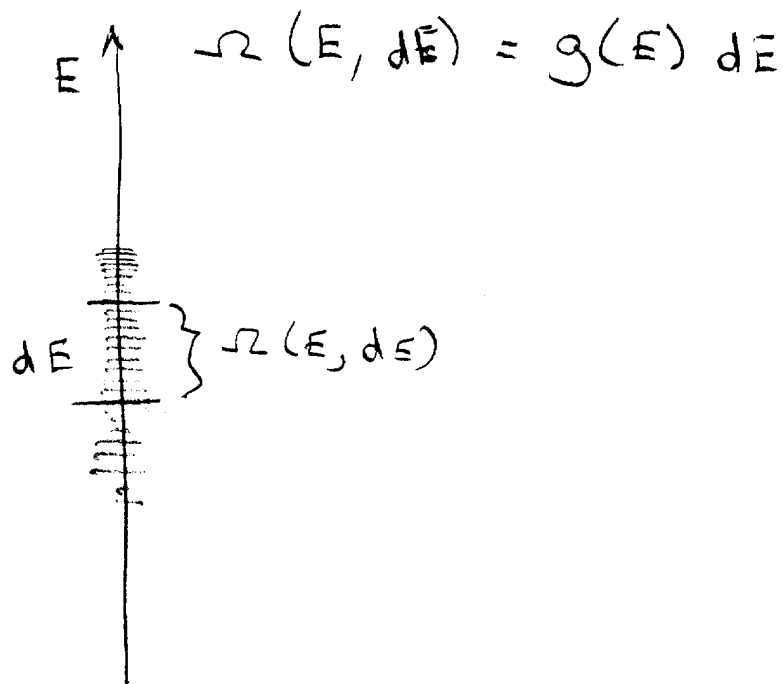
$$P_1 = \frac{\Omega_1}{\Omega_0} = \frac{2}{4}$$

## C The Spacing of States

The number of possible states for, say, molecules in a room so that  $\frac{1}{2}$  are in front &  $\frac{1}{2}$  are in back is  $\sim 10^{10^{28}}$ !

Most of these states have molecules  $\sim$  same energy. Spacing in energy  $\sim 10^{-24}$  eV (quantum levels)

We can describe this using the density of states:



## D. Density of States and the Internal Energy

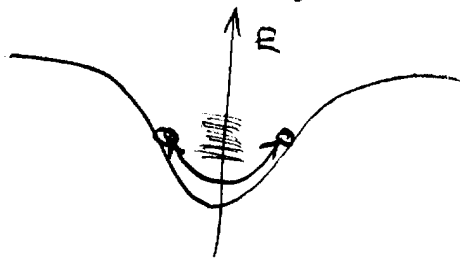
Let  $\Omega_i = \#$  of states in the  $i^{\text{th}}$  degree of freedom  
(say vibration in x direction)  
etc.

For each of these states there are  $\Omega_j$  states in the  
 $j^{\text{th}}$  degree of freedom, etc.

So in total there are

$$\begin{aligned}\Omega_0 &= \Omega_1 \cdot \Omega_2 \cdot \Omega_3 \cdots \Omega_R \quad \text{states} \\ &= \prod_{i=1}^R \Omega_i \quad \uparrow \# \text{ of Dof.}\end{aligned}$$

Now as  $E$  in a system  $\uparrow$ , there are more available states



$$\Omega_i(E) \propto E^\chi \leftarrow \text{some positive number. (typically } \frac{1}{2} \text{)}$$

$$\therefore \Omega_0 = \Omega_1 \cdot \Omega_2 \cdot \Omega_3 \cdots \propto (E^\chi)^R = E^{\chi R}$$

$$\text{Now } R \sim 10^{22} \rightarrow 10^{28}$$

$$\therefore \Omega_0 \uparrow \text{ very rapidly as } E \uparrow$$

What is  $\chi$ ?

we had (see 5-1)

$$E = b q^2, \text{ eg } \frac{1}{2} m v^2, \frac{1}{2} k x^2, \dots$$

So max energy in a given D.O.F.,  $= b q_{\max}^2$

$\therefore$  # of accessible states,  $\Omega_i \propto q_{\max} \propto E^{1/2}$   
ie  $\chi = \frac{1}{2}$ .

$$\therefore \Omega_0(E) = \prod_{i=1}^R \Omega_i = \text{const} \times E^{R/2}$$

$$\dagger \Omega_0(E, dE) = \underbrace{\text{constant} E^{R/2}}_{g(E)} dE$$

$\uparrow$  density of states

Example

$$\text{if } R = 2 \times 10^{24} \Rightarrow g(E) \propto E^{10^{24}} !$$

as  $E \uparrow$   $g \uparrow$  very rapidly.

Example Room full of air  $10^{28}$  molecules  
 5 DOF for each molecule  
 (3 translational, rotational)

T increases by  $1^\circ\text{C}$ .

By what factor does # states increase?

Sol'n

$$T \sim 300\text{K}$$

$$E \sim kT \text{ so } \frac{E_2}{E_1} \sim \frac{301}{300} = 1.0033$$

$$\therefore \frac{\Omega_0^{301}}{\Omega_0^{300}} = \left(\frac{E_2}{E_1}\right)^{R/2} = (1.0033)^{\frac{5}{2} \times 10^{28}}$$

$$= 10^{3.6 \times 10^{26}}$$

Huge increase in # states when T  $\uparrow$  by 1 degree.

$$\underline{\Omega_0^{301} = \Omega_0^{300} \times 10^{3.6 \times 10^{26}}}$$

!