

## Chapter 22 Heat Engines, Entropy and the Second Law of Thermodynamics

First law states that energy is conserved.

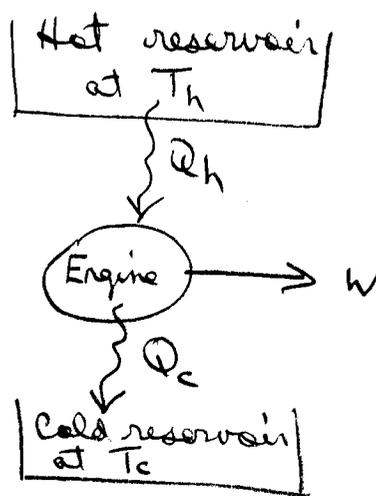
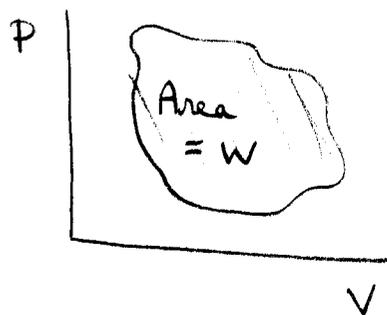
Second law establishes which processes can and cannot occur.

### 22.1 Heat Engines and the Second Law of Thermodynamics

Heat engine:

- Working substance

- cyclic process  $\Rightarrow \Delta U = 0 \Rightarrow W = Q_h - Q_c$



Thermal Efficiency,  $e \equiv \frac{W}{Q_h} = \frac{Q_h - Q_c}{Q_h} = 1 - \frac{Q_c}{Q_h}$

Typically,  $e \sim 20 \rightarrow 40\%$  in engines.

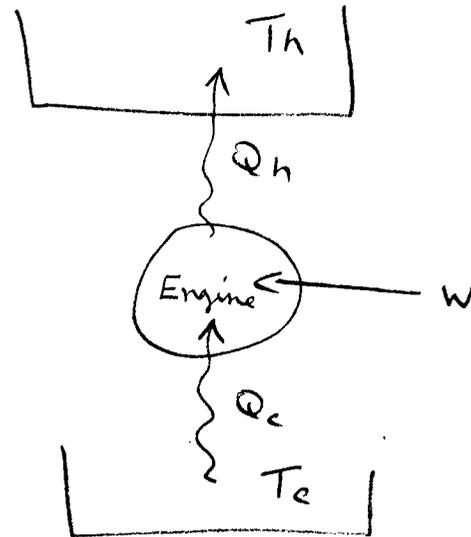
2<sup>nd</sup> Law: It is impossible to construct a heat engine that, operating in a cycle, produces no other effect than the absorption of thermal energy from a reservoir and the performance of an equal amount of work. (Kelvin-Planck)

- in short:  $Q_c \neq 0$

## Refrigerators and Heat Pumps

These are heat engines in reverse.

Must do work,  $W$ , to cause thermal energy to flow from cold reservoir to hot reservoir. (Clausius)

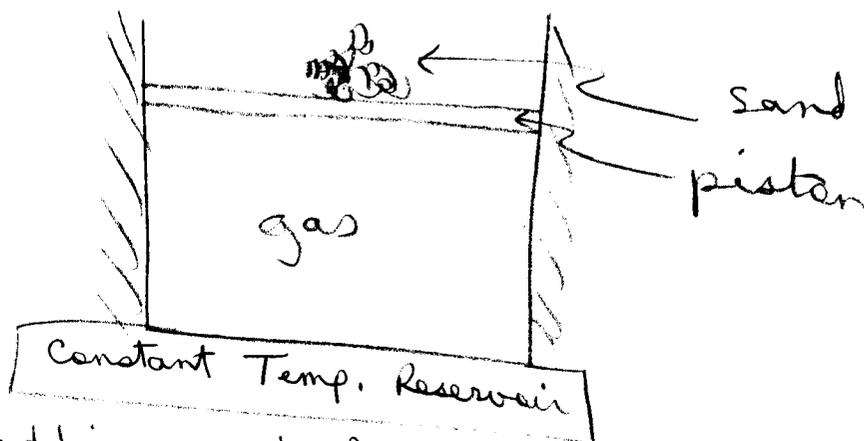


## 22.2 Reversible and Irreversible Processes

Reversible Process: one that can be performed so that, at its conclusion, both the system and its surroundings have been returned to their exact initial conditions.

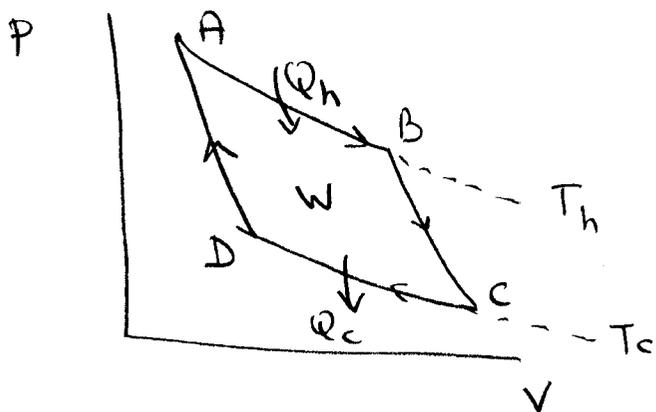
Irreversible Process: one that does not satisfy the above.

All real processes are irreversible but in some cases we can approximate a reversible process by minimizing dissipation effects such as friction.



slowly adding grains of sand causes a reversible compression of the gas - i.e. the system looks reversible but the surroundings are not.

## 22.3 The Carnot Engine

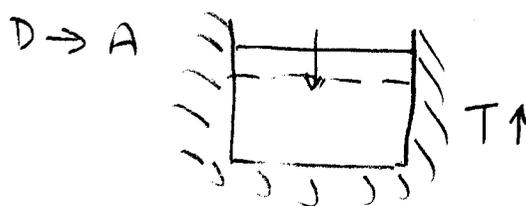
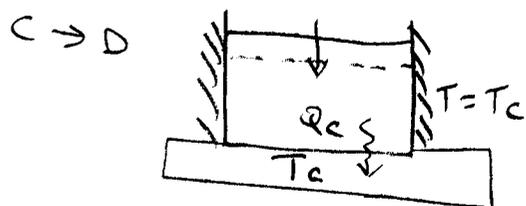
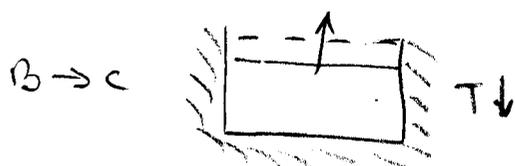
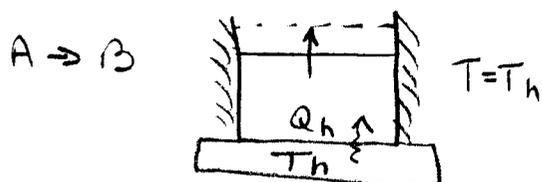


A  $\rightarrow$  B Isothermal expansion at  $T_h$   
 $Q = Q_h$  (= heat added)

B  $\rightarrow$  C Adiabatic expansion  
 $Q = 0$

C  $\rightarrow$  D Isothermal compression at  $T_c$   
 $Q = Q_c$  (= heat rejected)

D  $\rightarrow$  A Adiabatic compression  
 $Q = 0$



A → BT constant  $\therefore \Delta U_{\text{gas}} = 0$ 

$$\therefore Q_h = W_{AB} = nRT_h \ln(V_B/V_A)$$

C → D

Similar to A → B

$$\therefore Q_c = -W_{CD} = nRT_c \ln(V_c/V_D)$$

B → C

Adiabatic, quasi-static expansion.

$$\therefore PV^\gamma = \text{constant}$$

$$\& PV = nRT$$

$$\Rightarrow TV^{\gamma-1} = \text{constant}$$

$$\therefore T_h V_B^{\gamma-1} = T_c V_c^{\gamma-1}$$

D → A

Similar to B → C

$$\therefore T_h V_A^{\gamma-1} = T_c V_D^{\gamma-1}$$

$$\therefore \left(\frac{V_D}{V_A}\right)^{\gamma-1} = \left(\frac{V_c}{V_B}\right)^{\gamma-1} \Rightarrow \frac{V_B}{V_A} = \frac{V_c}{V_D}$$

$$\therefore \frac{Q_c}{Q_h} = \frac{nRT_c \ln(V_c/V_D)}{nRT_h \ln(V_B/V_A)} = \frac{T_c}{T_h}$$

$$\therefore e = 1 - \frac{Q_c}{Q_h} = 1 - \frac{T_c}{T_h} \leftarrow \text{for Carnot Cycle}$$

from Serway Chapter 21:

$$du = nC_v dT = \underbrace{-PdV}_{-W}$$

but  $PV = nRT$  if ideal gas

$$\therefore PdV + VdP = nRdT \leftarrow$$

$$\therefore PdV + VdP = -\frac{R}{C_v} PdV$$

$$\text{Now } R = C_p - C_v$$

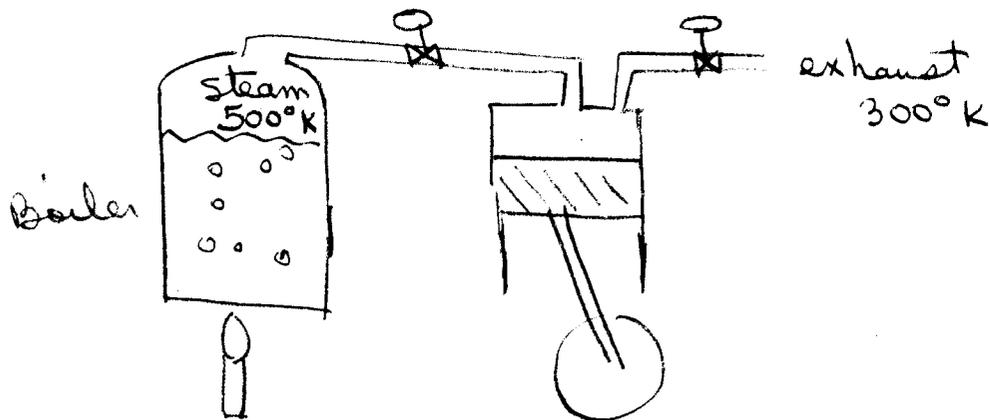
$$\therefore \frac{dP}{P} + \gamma \frac{dV}{V} = 0$$

$$\text{where } \gamma = C_p/C_v$$

$$\Rightarrow \ln P + \gamma \ln V = \text{constant}$$

$$\therefore PV^\gamma = \text{constant.}$$

### Example: The Steam Engine



Find maximum efficiency.

$$e_{\text{Carnot}} = 1 - \frac{T_c}{T_h} = 1 - \frac{300\text{K}}{500\text{K}} = 0.4 \text{ or } \underline{40\%}$$

Find  $W_{\text{max}}$  if  $Q_h = 200 \text{ J}$ .

$$e = 1 - \frac{Q_c}{Q_h} = \frac{W}{Q_h} = 0.4 \quad \therefore \quad W = 0.4 \times 200 \text{ J} = \underline{80 \text{ J}}$$

### Example: The Carnot Efficiency

Gas engine,  $e = 30\%$ ,  $T_c = 300\text{K}$ ,  $T_h = ?$   
(assume Carnot)

$$e = 0.3 = 1 - \frac{T_c}{T_h} \quad \therefore \quad T_h = \frac{T_c}{1-e} = \frac{300}{.7} = \underline{\underline{492 \text{ K}}}$$

If  $Q_h = 837 \text{ J}$  find  $W$ .

$$W = e Q_h = 0.3 \times 837 \text{ J} = \underline{\underline{251 \text{ J}}}$$

## 22.4 The Absolute Temperature Scale

We can use the Carnot cycle as a means of measuring temperature.

Choose a reference temperature: 273.16 K  
 (triple point of water, i.e. the temp. at which water in vapour, liquid and ice coexist -  
 $T = 0.01^\circ\text{C}$ ,  $P = 4.58 \text{ mm. Hg}$ )

For Carnot cycle we showed that

$$\frac{T_c}{T_h} = \frac{Q_c}{Q_h}$$

- Run a substance through the Carnot cycle wherein  $T_c = 273.16 \text{ K}$ .
- Measure  $Q_c + Q_h$ .
- Infer  $T_h$  from  $T_h = (273.16 \text{ K}) \times \frac{Q_h}{Q_c}$

## 22.5 The Gasoline Engine

### Otto cycle

$O \rightarrow A$  Intake stroke  
Vol increases from  
 $V_2 \rightarrow V_1$

$A \rightarrow B$  Compression stroke

Adiabatic

$V_1 \rightarrow V_2$ ,  $T_A \rightarrow T_B$  (T increase)

Work of compression = area under AB

$B \rightarrow C$  Combustion,  $Q_h$  added

$V \sim$  constant,  $P \uparrow$ ,  $T \uparrow$ ,  $W = 0$

$C \rightarrow D$  Power stroke

Adiabatic expansion,  $V_2 \rightarrow V_1$

$T \downarrow$ ,  $T_c \rightarrow T_0$

Work of expansion = area under CD

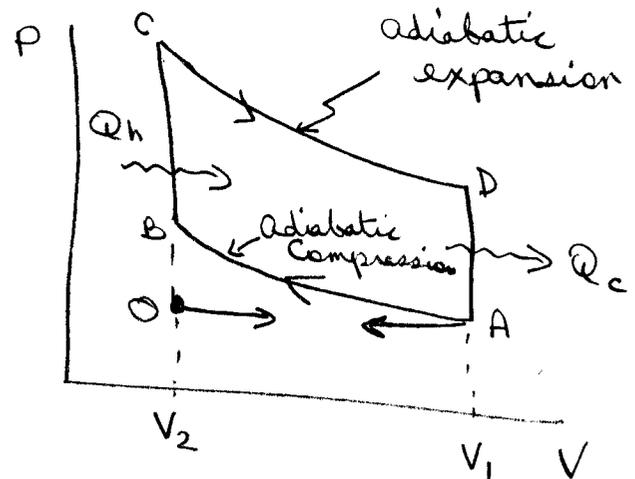
$D \rightarrow A$  Exhaust valves open

$P \downarrow$ ,  $V$  constant,  $W = 0$ ,  $Q_c$  removed

$A \rightarrow O$  Exhaust stroke

Residual gas removal

$V \downarrow$ ,  $V_1 \rightarrow V_2$



## Example: Efficiency of the Otto Cycle

Assume ideal gas.  $\Delta U = 0$  for a cycle.

$$\therefore W = Q_h - Q_c$$

$$Q_h = n C_v (T_c - T_b) \quad (\text{constant volume heat addition})$$

$$Q_c = n C_v (T_d - T_a) \quad (\text{constant volume heat removal})$$

$$\therefore e = \frac{W}{Q_h} = 1 - \frac{Q_c}{Q_h} = 1 - \frac{(T_d - T_a)}{(T_c - T_b)}$$

Now since  $A \rightarrow B$  +  $C \rightarrow D$  are adiabatic:

$$T_A V_1^{\gamma-1} = T_B V_2^{\gamma-1} \quad , \quad \gamma = C_p/C_v \approx 1.4$$

and  $T_D V_1^{\gamma-1} = T_C V_2^{\gamma-1}$

$$\begin{aligned} \therefore \frac{T_D - T_A}{T_c - T_b} &= \frac{T_C (V_2/V_1)^{\gamma-1} - T_B (V_2/V_1)^{\gamma-1}}{T_D - T_B} \\ &= (V_2/V_1)^{\gamma-1} \end{aligned}$$

$$\therefore e = 1 - \frac{1}{(V_1/V_2)^{\gamma-1}}$$

$V_1/V_2$  is called the compression ratio.

If  $V_1/V_2 = 8$  (typical)  
then  $e = 56\%$

$e \uparrow$  as  $V_1/V_2 \uparrow$

Actual is  $\sim 20\%$  due to friction, heat losses + incomplete combustion.

In terms of temperatures:

$$T_A V_1^{\gamma-1} = T_B V_2^{\gamma-1}$$

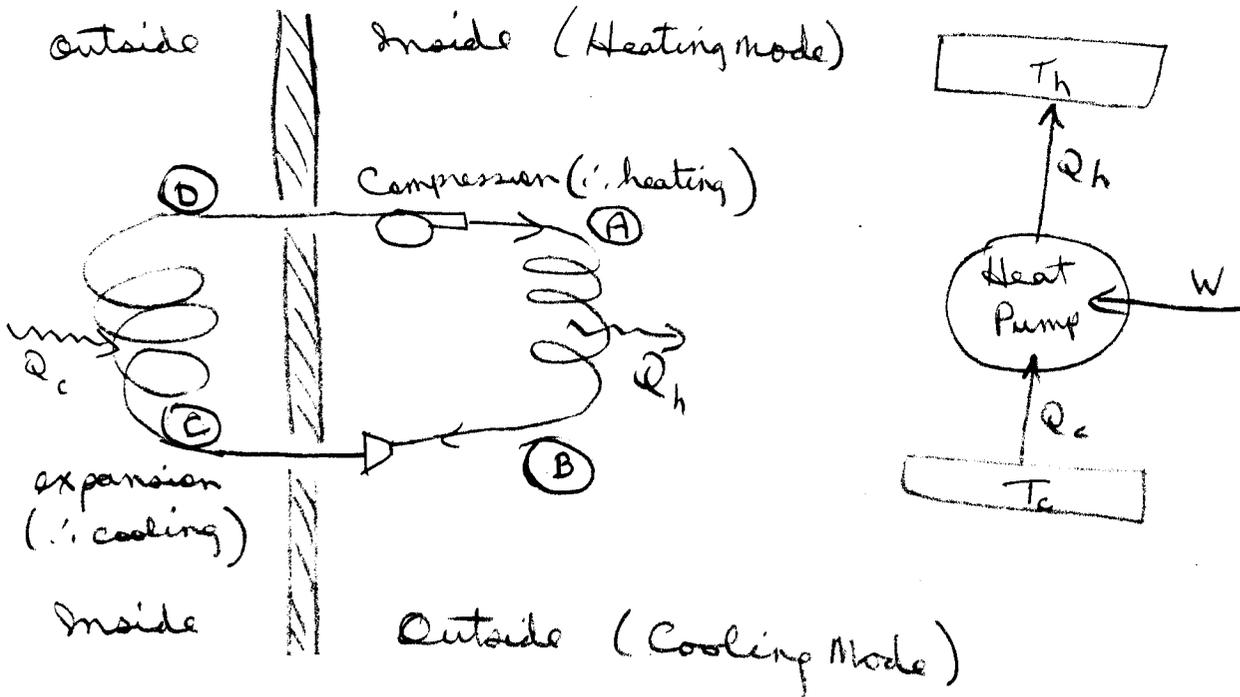
$$\therefore \frac{T_A}{T_B} = \left(\frac{V_2}{V_1}\right)^{\gamma-1} = \frac{T_D}{T_C}$$

$$\therefore e = 1 - \frac{T_A}{T_B} = 1 - \frac{T_D}{T_C}$$

Compare to a Carnot cycle:

$$e_{\text{Carnot}} = 1 - \frac{T_{\text{lowest}}}{T_{\text{highest}}} = 1 - \frac{T_A}{T_C} > e_{\text{Otto}}$$

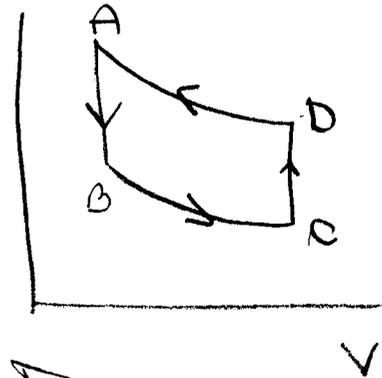
# 22.6 Heat Pumps and Refrigerators



Carnot cycle in reverse

Coefficient of performance for heat pump

$$\equiv \text{COP} \equiv \frac{Q_h}{W} = \frac{Q_h}{Q_h - Q_c}$$



$$\begin{aligned} \text{COP}_{\text{carnot}} &= \text{COP}_{\text{maximum}} \\ &= \frac{T_h}{T_h - T_c} \end{aligned}$$

typically  $\sim 4$  when  $T_c \sim 25^\circ\text{C}$ .

$$\text{COP for refrigerator} \equiv \frac{Q_c}{W} \quad (\text{typically } \sim 5 \text{ or } 6)$$

$$\therefore \text{COP}_{\text{max}} (\text{refrigerator}) = \frac{T_c}{T_h - T_c}$$